High Resolution Pulsed Field Ionization Photoelectron Study of O_2 : Predissociation Lifetimes and High-n Rydberg Lifetimes Converging to $O_2^+(B^2\Sigma_\sigma^-, v^+=0)$ and 5)

C.-W. Hsu¹, M. Evans², S. Stimson², and C. Y. Ng²

¹Chemical Science Division, Ernest Orlando Lawrence Berkeley National Laboratory,
University of California, Berkeley, CA 94720, USA

²Ames Laboratory, USDOE and Department of Chemistry,
Iowa State University, Ames, IA 50011, USA

INTRODUCTION

The rotationally resolved photoelectron bands for $O_2^+(B^2\Sigma_g^-, v^+=0 \text{ and } 5)$ in the energy range of 20.27-20.95 eV have been measured using synchrotron based pulsed field ionization photoelectron techniques at an instrumental resolution of 5 cm⁻¹ (full-width-at-half-maximum). In addition to the determination of accurate ionization energies and rotational constants, we have also obtained the predissociative lifetimes of 0.9 ± 0.3 ps for $O_2^+(B^2\Sigma_g^-, v^+=0)$ and 0.50 ± 0.06 ps for $O_2^+(B^2\Sigma_g^-, v^+=0)$. The (nominal) effective lifetimes for high-n Rydberg states converging to $O_2^+(B^2\Sigma_g^-, v^+=0)$ and 5) are measured to be $\approx 0.4~\mu s$, which are significantly shorter than those of $\approx 1.9~\mu s$ observed for $O_2^+(b^4\Sigma_g^-, v^+=0.5)$. The shorter (nominal) effective lifetimes for high-n Rydberg states converging to $O_2^+(B^2\Sigma_g^-, v^+=0)$ and 5) are attributed to the higher kinetic energy releases (or velocities) of $O_2^+(B^2\Sigma_g^-, v^+=0)$ and 5) ion cores.

EXPERIMENT

In the supersonic beam experiment, a continuous O_2 beam was produced by supersonic expansion of pure O_2 through a stainless steel nozzle (diameter=0.127 mm, T=298 K) at a stagnation pressure of 760 Torr. As shown below, the simulation of PFI-PE spectra obtained using a supersonically cooled O_2 sample indicates that the rotational temperature of O_2 achieved is ≈ 35 K. The O_2 sample intersects the monochromatic VUV light beam 7 cm downstream in the photoionization/photoexcitation (PI/PEX) region. We have also performed PFI-PE measurements using an O_2

effusive beam, which was introduced into the PI/PEX region by a metal orifice (diameter= 0.5 mm) at room temperature and a distance of 0.5 cm from the PI/PEX region.

In the present experiment,¹ the nominal dc electrostatic field at the PI/PEX region was zero by setting the repeller plates at the same potential before the application of the Stark electric field pulse. A pulsed electric field (height=1.1V/cm, width=40 ns, delayed by 20 ns with respect to the beginning of the 60 ns synchrotron dark gap) was applied to the repeller at the PI/PEX region every one (or two, or three) ring period. The pulsed electric field was used to ionize the high-n Rydberg states and extract the PFI-PEs toward the detector. As demonstrated in the Ne⁺(²P_{3/2}) PFI-PE band shown in Fig. 1, the PFI-PE resolution achieved here was 0.63±0.05 meV or 5.0±0.4 cm⁻¹ (FWHM) at 21.5648 eV.

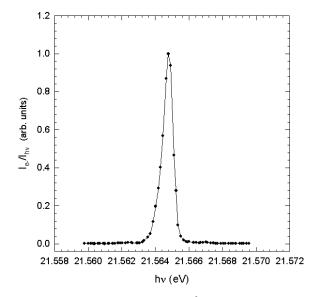


Figure 1. PFI-PE spectra of $Ne^+(^2P_{3/2})$ obtained using monochromator entrance/exit slits of 50/50 μm . The PFI-PE resolution achieved was 0.63 \pm 0.05 meV.

RESULTS

Figures 2(a) and 2(b) show the PFI-PE bands for $O_2^+(B^2\Sigma_g^-, v^+=0)$ (upper curves, open circles) in the photon energy range of 20.670-20.710 eV obtained using a supersonic O_2 beam and an effusive O_2 sample, respectively. The PFI-PE band for $O_2^+(B^2\Sigma_g^-, v^+=5)$ in the energy range of 20.905-20.945 eV observed using an effusive O_2 sample is depicted in Fig. 3 (upper curve, open circles). We note that the rotational features resolved for the $v^+=0$ band are narrower than those for the $v^+=5$ band.

The relative rotational intensities in individual vibrational bands have been simulated using the Buckingham-Orr-Sichel (BOS) model,² which was derived to predict rotational line strengths observed in one photon ionization of diatomic molecules. Due to nuclear spin statistics, the even levels of N" in $O_2(X^3\Pi_{\circ})$ and N⁺ in $O_2^+(B^2\Sigma_{\alpha})$ do not exist. Thus, only the rotational branches with $\Delta N (= N^+-N^") = \text{even}$ are possible. The observed rotational branches $\Delta N = -2$, 0, and +2 (O, Q, and S) branches, are marked in Figs 2(a), 2(b), and 3. For a $g \leftrightarrow g$ transition, the photoelectron angular momentum l must be odd. Thus, the partial waves for the ejected electron are restricted to the *l*=1 and 3 continuum states.

The simulation of the PFI-PE bands for v+=0 and 5 yields ionization energies (IEs) of $20.2982_5\pm0.0005$ and $20.9348_8\pm0.0005$ eV for the formation of $O_2^+(B^2\Sigma_g^-, v^+=0, N^+=1)$ and $O_2^+(B^2\Sigma_g^-, v^+=5, N^+=1)$ from $O_2(X^3\Sigma_g^-, v^*=0, N^*=1)$. The IEs are higher than the literature values by 2.3 and 8 meV, respectively. The simulation also yields rotational constants of 1.243 ± 0.002 cm-1 for the v+=0 state and 1.122 ± 0.002 cm-1 for the v+=5 state. These values allow the calculation of the equilibrium bond distances of re = 1.302 ± 0.001 and 1.370 ± 0.001 Å for the v+=0 and 5 states of $O2+(B2\Sigma_g^-)$, respectively. The latter values are greater than the r_e value of 1.2074 Å for

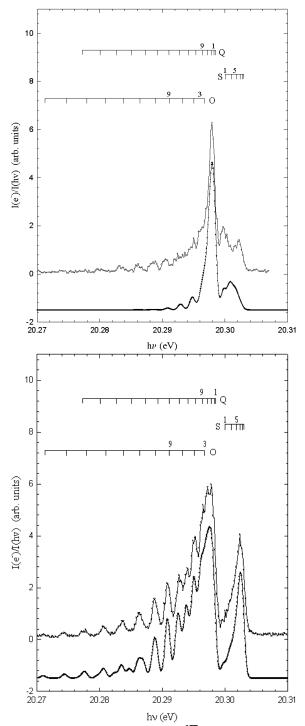


Figure 2. PFI-PE bands for $O_2^+(B^2\Sigma_g^-, v^+=0)$ (upper curves, open circles) obtained using (a) a supersonically cooled O2 sample and (b) an effusive O2 sample. The simulated spectra (lower curves, solid circles) were obtained using rotational temperatures of (a) 35 K and (b) 298 K.

 $O_2(X^3\Sigma_g^-, v"=0)$. This observation is consistent with the fact that the formation of the $O_2^+(B^2\Sigma_g^-)$ state involves the ejection of a bonding electron from the $3\sigma_g$ orbital.

The $O_2^+(B^2\Sigma_{\alpha}^-)$ state lies significantly above the first dissociation limit of $[O^{+}(^{4}S) + O(^{3}P)]$ at 18.73 eV and is known to be strongly predissociative. This, together with the fact that the energy of $O_2^+(B^2\Sigma_g^-)$ is more than 8 eV higher than that of the ground $O_2^+(X^2\Pi_g)$ state, makes spectroscopic and dynamical studies of $O_2^+(B^2\Sigma_g^-, v^+)$ difficult by employing the ion emission and common laser spectroscopic techniques. Although the $O_2^+(B^2\Sigma_g^-) \rightarrow O_2^+(A^2\Pi_u)$ transition is optically allowed, the emission from $O_2^+(B^2\Sigma_{\sigma})$ has not been observed. This suggests that the predissociative lifetimes $(\tau_d$'s) are significantly shorter than the radiative lifetimes for $O_2^+(B^2\Sigma_g^-, v^+)$. The predissociation dynamics of $O_2^+(B^2\Sigma_{s}^-, v^+)$ have been investigated extensively using photoelectron photoion coincidence (PEPICO) time-of-flight (TOF) techniques.³ In these experiments, the measurement of the fragment ion TOF spectrum is triggered by the detection of a threshold photoelectron (TPE). Since the TOF of a TPE to the

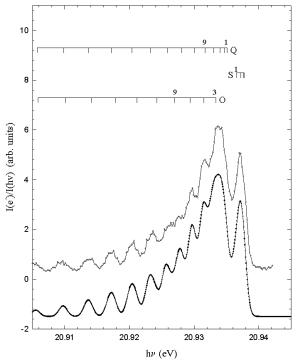


Figure 3. PFI-PE bands for $O_2^+(B^2\Sigma_g^-, v^+=5)$ (upper curves, open circles) obtained using an effusive O_2 sample. The simulated spectrum (lower curves, solid circles) was obtained using a rotational temperature of 298 K.

electron detector usually ranges from 10–100 ns, the traditional PEPICO-TOF technique is not applicable for the measurement of dissociative lifetimes shorter than 10 ns.

From the Gaussian linewidths used in the BOS simulation, we calculated the natural rotational linewidths for the respective v+=0 and 5 states to be 6.2±1.8 cm⁻¹ and 10.9±1.3 cm⁻¹. Since the τ_d 's are expected to be significantly shorter than the radiative lifetimes for $O_2^+(B^2\Sigma_g^-, v^+)$, the broadening of the rotational transitions is contributed predominantly to predissociation. Using these natural rotational linewidths we calculate the τ_d values of $(0.9\pm0.3)x10^{-12}~s$ for $O_2^+(B^2\Sigma_g^-, v^+=0)$ and $(0.50\pm0.06)x10^{-12}~s$ for $O_2^+(B^2\Sigma_g^-, v^+=5)$.

The τ values for high-n Rydberg states converging to $O_2+(B^2\Sigma_u^-, v^+=0 \text{ and } 5)$ are measured to be $\approx 0.4~\mu s$. In a similar study,4 we found that the τ values are $\approx 1.9~\mu s$ for high-n Rydberg states converging to the $O_2^+(b^4\Sigma_g^-, v^+=0-5)$ states, and are nearly independent of v^+ . Since the $O_2^+(b^4\Sigma_g^-, v^+=4 \text{ and } 5)$ states are known to be predissociative with τ_d values in the range of 0.01-4 ns, it is surprising that the τ values for high-n Rydberg states converging to these unstable ion cores are found to be nearly the same as those converging to the stable $O_2^+(b^4\Sigma_g^-, v^+=0-3)$ ion cores. For an electron in a sufficiently high-n and high-l Rydberg state with a near circular orbital, the Rydberg electron may not respond rapidly even when the core is dissociating. This would result in a substantially longer (autoionization) lifetime for the high-n Rydberg state compared to the τd value of the ion core. The latter conclusion should be valid if the kinetic energy for the departing $O^+ + O$ is sufficiently small. Thus, the (autoionization) lifetime of a high-n Rydberg state is effectively decoupled from the dissociative lifetime of the ion core.

Assuming that the autoionization and fluorescence lifetimes for a high-n Rydberg O_2 state, $O_2(n)$, are longer than the τ_d value of the O_2^+ ion core, a plausible decay mechanism for $O_2(n)$ is shown in reactions (1a) and (1b).

$$O_2(n) \rightarrow O(n') + O,$$
 (1a)
 $\rightarrow O^+ + e^- + O$ (1b)

If the relative velocity for the departing $O^+ + O$ fragment pair resulting from the predissociation of the O_2^+ ion core is not too high, the high-n Rydberg electron originally associated with O_2^+ in $O_2(n)$ may be guided by the Coulombic field to orbit around the departing O^+ ion forming a high-n' Rydberg O atom, O(n'). The relative velocities for the departing $O^+(^4S) + O(^3P)$ fragments from $O_2^+(B^2\Sigma_u^-, v^+=0)$ and $O_2^+(B^2\Sigma_u^-, v^+=5)$ are calculated to be 0.62×10^6 and 0.72×10^6 cm/s, respectively. The corresponding times required for O(n') fragments to move a distance of 0.15 cm, which defines the detection zone of the electron spectrometer, are 0.48 and 0.42 μ s. Since these values are comparable to the τ values of $\approx 0.4~\mu$ s observed for high-n Rydberg states converging to $O_2^+(B^2\Sigma_u^-, v^+=0)$ and 5), we conclude that the higher velocities (or kinetic energies) for O(n') formed in the dissociation reaction (1a) contribute to the short τ values for high-n Rydberg states converging to $O_2^+(B^2\Sigma_g^-, v^+=0)$ and 5). The τ values presented here must be considered as nominal values.

SUMMARY

We have obtained accurate spectroscopic constants and reliable τ_d values for $O_2^+(B^2\Sigma_g^-, v^+=0)$ and 5) using high resolution synchrotron based PFI-PE techniques. This method is directly applicable for τ_d measurements of other predissociative states of O_2^+ and predissociative states of other diatomic molecular ions in the inner-valence region. The results for τ measurements of $O_2^+(b^4\Sigma_g^-, B^2\Sigma_g^-)$ show that a high-n Rydberg molecular state with a predissociative molecular ion core is well defined only prior to the dissociation of the molecular ion core. Similar experiments have been performed on the $O_2^+(c^4\Sigma_u^-, v^+=0, 1)$ states and the results are soon to be published.

ACKNOWLEDGMENTS

C.Y.N. acknowledges helpful discussion with Prof. Tomas Baer.

REFERENCES

- 1. C.-W. Hsu, M. Evans. P. A. Heimann, and C. Y. Ng, Rev. Sci. Instrum., 68, 1694 (1997).
- 2. A. D. Buckingham, B. J. Orr, J. M. Sichel, Phil. Trans. Roy. Soc. Lond. A 268, 147 (1970).
- 3. T. Akahori, Y. Morioka, M. Watanabe, T. Hayaishi, K. Ito, and M. Nakamura, J. Phys. B 18, 2219 (1985).
- 4. C.-W. Hsu, M. Evans, S. Stimson, C. Y. Ng, and P. Heimann, Chem. Phys., accepted.

This work was supported by the Director, Office of Energy Research, Office of Basic Energy Sciences, Chemical Sciences Division of the U.S. Department of Energy under Contract No. DE-AC03-76SF00098 for the Lawrence Berkeley National Laboratory and Contract No. W-7405-Eng-82 for the Ames Laboratory.

Principal investigator: Prof. C.Y. Ng, Ames Laboratory, USDOE and Department of Chemistry, Iowa State University, Ames, IA 50011, USA. Email: cyng@ameslab.gov. Telephone: 515-294-4225